

## Research Article

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Feng-Xia Liang, Cai-Wang Ge, Teng-Fei Zhang, Wei-Jie Xie, Deng-Yue Zhang, Yi-Feng Zou, Kun-Zheng, and Lin-Bao Luo\*

# Plasmonic hollow gold nanoparticles induced high-performance $\text{Bi}_2\text{S}_3$ nanoribbon photodetector

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**Abstract:** A high performance hollow gold nanoparticles (HGNs) decorated one-dimensional (1-D)  $\text{Bi}_2\text{S}_3$  nanoribbon (NR) photodetector was fabricated for green light detection (560 nm). The single crystal 1-D  $\text{Bi}_2\text{S}_3$  NRs with growth orientation along [001] were synthesized by a simple solvothermal approach. Optoelectronic analysis reveals that the performance of the plasmonic photodetector was greatly enhanced after decoration with HGNs. For example, the responsivity increases from  $1.4 \times 10^2$  to  $1.09 \times 10^3 \text{ AW}^{-1}$ , the conductivity gain from  $2.68 \times 10^2$  to  $2.31 \times 10^3$ , and the detectivity from  $2.45 \times 10^{12}$  to  $2.78 \times 10^{13}$ , respectively. Such performance enhancement was attributed to the localized surface plasmon resonance (LSPR) effect caused by the HGNs according to both experiment and theoretical simulation. This study is believed to open up new opportunities for managing light and enhancing the device performance of other 1-D semiconductor nanostructures based optoelectronic devices and systems.

**Keywords:** Green light detection, localized surface plasmon resonance, responsivity, finite element method, plasmonic nanoparticles

## 1 Introduction

Currently, noble metal nanoparticles (e.g. Au, Ag, and Pt) have been extensively studied in a variety of field due to their localized surface plasmon resonance (LSPR) effect,

which is actually caused by the collective oscillations of the excited free electrons upon light illumination with a certain wavelength [1], [2], [3]. LSPR could couple with the incident light, leading to more light trapped nearby semiconductor in plasmonic optoelectronic devices [4], [5]. This mode of manipulating light through engineering noble metal nanoparticles on semiconductors has been exploited in a number of applications, such as photocatalysis [6], [7], plasmonic nanoantenna [8], [9], solar cells [10], [11], light emitting diodes [12], and photodetectors [13, 14]. Among these light-harvesting devices, high performance photodetectors enabled by LSPR are receiving special research interest for their great potential in environment monitoring, biological analysis, industrial control, space-to-space communications and so on [15], [16], [17]. For instance, Sakurai *et al* developed a high performance Au nanoparticles@ZnO nanowire photodetectors. It was observed that the dark current decrease by 100 times and the on/off ratio was considerably increased after decoration of plasmonic Au nanoparticles [18]. Moreover, Hu's group observed that a large enhancement of photocurrent response is obtained by decorating multi-layer  $\text{MoS}_2$  with plasmonic gold nanoparticles array. Due to LSPR, the photocurrent of  $\text{MoS}_2$  phototransistors exhibited a threefold enhancement after depositing periodic Au nanoparticles arrays [19].

One dimensional (1-D) semiconductor nanostructures have demonstrated great potential in wide ranging device applications [20], [21], [22], [23]. Bismuth sulfide ( $\text{Bi}_2\text{S}_3$ ), as a typical semiconductor with a direct bandgap of 1.3 eV and a high absorption coefficient, has attracted great attention in hydrogen storage, electronic devices, and optoelectronic devices [24], [25], [26]. For instance, Konstanatos *et al.* reported a solution-processed  $\text{Bi}_2\text{S}_3$  nanocrystalline photodetector with the responsivity of about 40 A/W and a high detectivity on the order of  $10^{11}$  Jones [21]. Li *et al.* fabricated a visible light photodetector based on hierarchical  $\text{Bi}_2\text{S}_3$  nanostructures. The photodetector showed stable photoresponse with a fast response

\*Corresponding Author: Lin-Bao Luo: School of Electronic Sciences and Applied Physics, Hefei University of Technology, Hefei, Anhui 230009, P. R. China, E-mail: luolb@hfut.edu.cn

Feng-Xia Liang, Deng-Yue Zhang: School of Materials Science and Engineering, Hefei University of Technology, Hefei, Anhui 230009, P. R. China

Cai-Wang Ge, Teng-Fei Zhang, Wei-Jie Xie, Yi-Feng Zou, Kun-Zheng: School of Electronic Sciences and Applied Physics, Hefei University of Technology, Hefei, Anhui 230009, P. R. China

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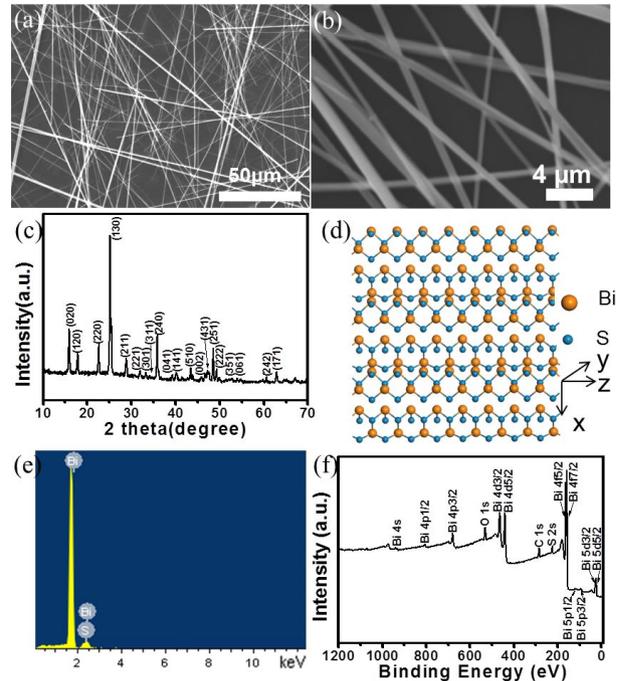
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time of 50 ms [27]. Despite these efforts, however, the device performance of the majority of the 1-D  $\text{Bi}_2\text{S}_3$  nanostructure based photodetector ever reported was limited by relatively low photoresponsivity and external quantum efficiency, which is related to the extremely small absorption cross-section of the  $\text{Bi}_2\text{S}_3$  [28], [29]. Enlightened by the obvious positive role the plasmonic nanoparticles played in optoelectronic device, we herein reported a hollow gold nanoparticles (HGNs) induced high-performance  $\text{Bi}_2\text{S}_3$  NRs photodetectors. It was found that the intensity of the light absorption increased after the NRs surface was modified with HGNs. Furthermore, an improved device performance with enhanced responsivity, detectivity and photoconductive gain was achieved. Theoretical simulation based finite element method (FEM) shows that the enhancement in devices performance is due to the strong LSPR effect of the HGNs.

## 2 Structural analysis of the $\text{Bi}_2\text{S}_3$ NRs

The 1-D  $\text{Bi}_2\text{S}_3$  nanomaterial was synthesized by a modified solvothermal approach [30]. Both Figure 1(a) and b show the typical field-emission scanning electron microscope (FESEM) images of the as-synthesized  $\text{Bi}_2\text{S}_3$  nanostructures. It is obvious that the product is composed of fiber-like nanostructures with aspect ratio (namely the length/diameter ratio) as high as 1000. The nanostructures are up to hundreds of micrometers in length and displayed high purity without appreciable contaminants, which is highly favorable for subsequent device construction. Figure 1(c) depicts the X-ray diffraction (XRD) pattern acquired from the product. All the strong diffraction peak in the range from  $15\text{--}70^\circ$  can be readily ascribed to orthorhombic lattices of  $\text{Bi}_2\text{S}_3$  (JCPDS Card number: 17-320), in agreement with literature result [31]. Figure 2d shows the perspective view of the orthorhombic  $\text{Bi}_2\text{S}_3$ , in which the orange and cyan balls denote the bismuth and sulfur atoms, respectively. Further energy dispersive spectrometer (EDS) analysis shown in Figure 1e revealed the presence of B and S elements with atomic ratio of 1.9:3.1, in consistency with the stoichiometric ratio of  $\text{Bi}_2\text{S}_3$  and the X-ray photoelectron spectroscopy analysis (Figure 1f).

The microstructure of the as-prepared  $\text{Bi}_2\text{S}_3$  NRs was then analyzed by both transmission electron microscopy (TEM) and high-resolution TEM. Figure 2a shows a typical TEM image of an individual  $\text{Bi}_2\text{S}_3$  NR, in which the NR has a very smooth surface. The HRTEM image in Figure 2b exhibits a lattice-spacing of 0.38 nm, which corre-

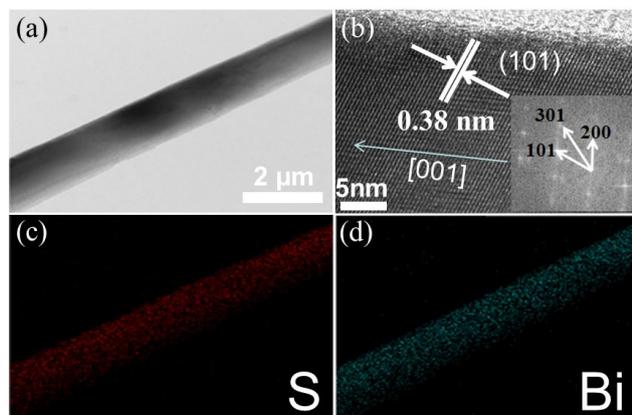


**Fig. 1.** (a) Representative SEM image of the  $\text{Bi}_2\text{S}_3$  NRs. (b) SEM image of the  $\text{Bi}_2\text{S}_3$  NRs at large magnification. (c) XRD pattern of the  $\text{Bi}_2\text{S}_3$  NRs. (d) Schematic illustration of the crystal structure of the  $\text{Bi}_2\text{S}_3$ . (e) EDX spectrum of  $\text{Bi}_2\text{S}_3$  NRs. (f) XPS spectrum analysis of the  $\text{Bi}_2\text{S}_3$  NRs.

sponds to the  $d$ -spacing of (101) plane. Further analysis of the HRTEM image as well as the selected area electron diffraction (SAED) pattern reveals a preferential growth direction along [001], which agrees with previous report [31]. The EDS elemental mapping shows that the constituting elements (S and Bi atoms) were uniform distributed in the NR (Figure 2c and d).

## 3 Results and discussion

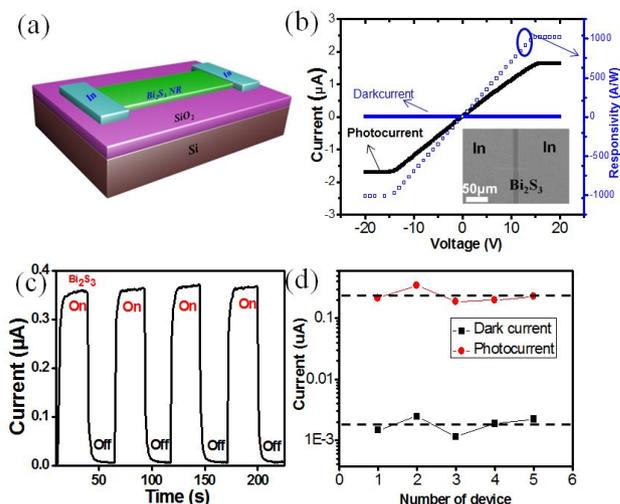
Next, nano-photodetector based on individual  $\text{Bi}_2\text{S}_3$  NR as shown in Figure 3(a) was fabricated to study its potential for device application. Figure 3(b) plots the typical  $I - V$  curves of an individual NR based photodetectors in dark and under 560 nm light illumination. It can be found that the  $I - V$  curves are linear at both conditions, suggesting that good contact (*Ohmic* contact) was formed at the In electrode/ $\text{Bi}_2\text{S}_3$  NRs interface. When 560 nm light was illuminated on the photodetector, the current increased abruptly from  $1.46 \times 10^{-3} \mu\text{A}$  (corresponding to voltage of 2 V) to  $0.215 \mu\text{A}$ , yielding an on/off ratio (i.e. photocurrent/dark current) of 147. Interestingly, with the increase of bias voltage, the photocurrent was observed to saturate



**Fig. 2.** (a) TEM image of an individual Bi<sub>2</sub>S<sub>3</sub> NR. (b) Corresponding HRTEM image. (c) and (d) elemental mapping of both S and Bi.

when the voltage is as high as  $\sim 14$  V. It should be noted that the above photoresponse is highly reproducible when the illumination was repeatedly switched on and off. As shown in Figure 3(c), the photocurrent quickly increases to the maximum value and remains at a steady state upon light illumination. This process could be repeated with the light on and off for several cycles, indicative of good reproducibility of the photodetector device. To obtain a statistical significance, the dark current and photocurrent of 5 representative devices were plotted (Figure 3d). It is noted that all the devices exhibit obvious sensitivity to light illumination, with two lines denoting the average  $I_{\text{dark}}$  and  $I_{\text{light}}$  of  $1.83 \times 10^{-3}$ , and  $2.36 \times 10^{-1}$   $\mu\text{A}$ , respectively. The large photocurrent is attributed to the photoconductive properties of Bi<sub>2</sub>S<sub>3</sub> NRs upon 560 nm light illumination, whose energy is larger than the bandgap of the Bi<sub>2</sub>S<sub>3</sub> NRs [32].

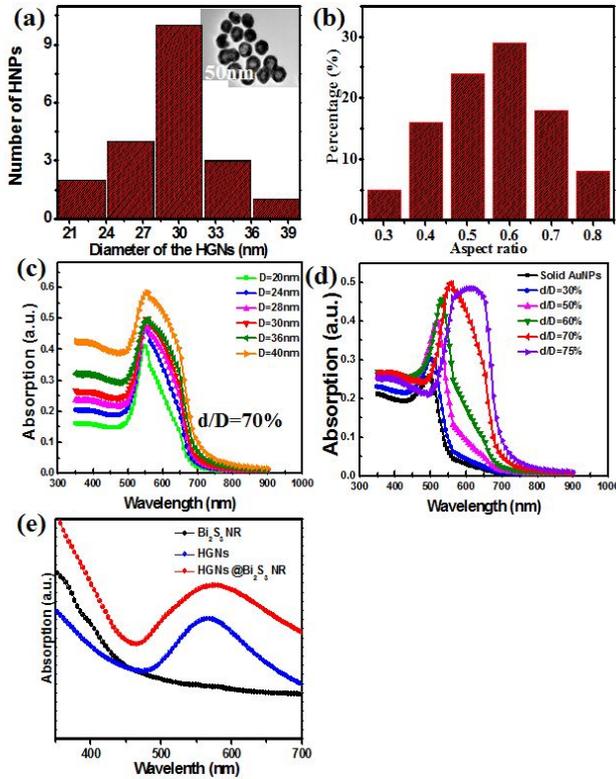
The inset of Figure 4(a) shows a typical TEM image of HGNS which were used to enhance the device performance of 1-D Bi<sub>2</sub>S<sub>3</sub> NR device. In this study, the HGNS were synthesized by a sacrificial galvanic replacement reaction between Co nanoparticles and aqueous HAuCl<sub>4</sub> [33], [34]. According to TEM analysis, their diameters are in the range from 20–40 nm, with average diameter of 30 nm (Figure 4b). What is more, the aspect ratio (i.e.: the inner-to-outer radius or  $d_{\text{in}}/d_{\text{outer}}$ ) ranges from 20–80%, with a mean value of 55%. Theoretical simulation based on finite element method (FEM) reveals that for all HGNS with varied diameters (the aspect ratio is set to be 70%), the absorption spectrum exhibits strong absorption at 560 nm, which can be attributed to the LSPR of the HGNS. Moreover, the absorption intensity increases with the increasing diameter. This finding is understandable as large metal HGNS are capable of increasing the optical path length inside the nanostructures. Figure 4(d) plots the theoretical



**Fig. 3.** (a) Schematic illustration of the Bi<sub>2</sub>S<sub>3</sub> NR device. (b)  $I - V$  curves and responsivity of Bi<sub>2</sub>S<sub>3</sub> NR in dark and under 560 nm light illumination, the inset shows an SEM image of the device. (c) Photoresponsive characteristics of the Bi<sub>2</sub>S<sub>3</sub> NR device. (d) Summary of both dark current and photocurrent of five representative individual Bi<sub>2</sub>S<sub>3</sub> NR based device without HGNS decoration.

simulation of absorption of solid gold nanoparticles (SGN) and HGNS with aspect ratio ranging from 30–70%. Among the various HGNS with different diameters, SGNs show the weakest light absorption [35]. Meanwhile, when the aspect ratio increases gradually from 0 to 30, 50, 60, 70, and 75%, the wavelength of LSPR band is observed to red-shift from 500 to 505, 517, 531, 554, and 613 nm. In fact, the LSPR of the HGNS was experimentally confirmed by UV-vis absorption spectra of HGNS, Bi<sub>2</sub>S<sub>3</sub> NRs and HGNS@Bi<sub>2</sub>S<sub>3</sub> NRs, as shown in Figure 4(e). The absorption of pure Bi<sub>2</sub>S<sub>3</sub> NRs is very weak at  $\sim 570$  nm. However, upon surface modification, the HGNS@Bi<sub>2</sub>S<sub>3</sub> NRs heterojunction exhibits strong absorption at 570 nm, which corresponds to the LSPR band of HGNS.

The proof-of-concept plasmonic photodetector was fabricated by modifying HGNS on Bi<sub>2</sub>S<sub>3</sub> NR surface via a simple chemical bonding method. Figure 5(a) illustrates the process to functionalize HGNS with ethylene glycol bis(3-mercaptopropionate), whose thiol groups (–SH) can bridge both Bi and Au atoms through chemical bonding. Optoelectronic study reveals that after HGNS decoration, the Bi<sub>2</sub>S<sub>3</sub> NR device displays apparently enhanced sensitivity to light illumination. The photocurrent at bias voltage 2 V is about 1.66  $\mu\text{A}$ , which is at least 6 times larger than that of device without HGNS [Figure 5(b)]. In addition, similar saturation in photocurrent was observed when the bias voltage was further increased. Our photoresponse study further suggests that



**Fig. 4.** (a) Statistical distribution of the HGNs diameter. (b) The inner-to-outer radius aspect ratio of the HGNs. (c) Theoretical simulation of the HGNs with different diameters, the diameters are from 20 to 24, 28, 30, 36, and 40 nm. In addition, the inner-to-outer radius aspect ratios of all HGNs are set to be 70%. (d) Theoretical simulation of absorption of HGNs, the HGNs have inner-to-outer radius ratios ranging from 0 to 75%, the diameters of all nanoparticles (including both solid AuNPs and HGNs) are set to be 30 nm. (e) Experimental absorption spectra of Bi<sub>2</sub>S<sub>3</sub> NRs, HGNs and HGNs @ Bi<sub>2</sub>S<sub>3</sub> NRs.

the HGNs@Bi<sub>2</sub>S<sub>3</sub> NR device is able to detect the illumination which was repeatedly switched on and off with good reproducibility [Figure 5(c)]. The dark current and photocurrent are  $1.87 \times 10^{-3}$  and  $1.61 \mu\text{A}$ , respectively. Figure 5d lists the dark current and photocurrent of 5 representative HGNs@Bi<sub>2</sub>S<sub>3</sub> NR devices. Compared with the device without HGNs (Figure 3d), the average photocurrent of HGNs@Bi<sub>2</sub>S<sub>3</sub> NR increased from  $2.36 \times 10^{-1}$  to  $1.93 \mu\text{A}$ , while the dark current keeps nearly unchanged ( $I_{\text{dark, Bi}_2\text{S}_3}/I_{\text{dark, HGNs@Bi}_2\text{S}_3}: 1.81 \times 10^{-3} \mu\text{A}/1.83 \times 10^{-3} \mu\text{A}$ ), leading to considerable increase in On/Off ratio (from 147 to 1066).

In order to quantitatively evaluate the performance of the present and other Bi<sub>2</sub>S<sub>3</sub> nanostructures based devices, several key parameters including responsivity ( $R$ ), conductivity gain ( $G$ ), and detectivity ( $D$ ) are calculated using

the following equations:[36], [37]

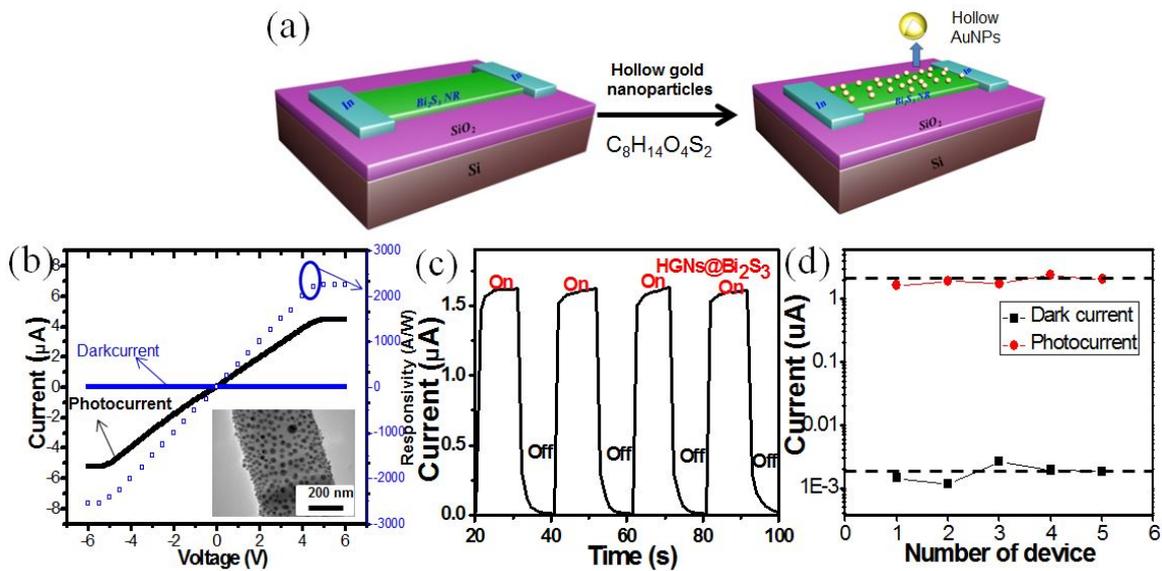
$$R(AW^{-1}) = \left( \frac{I_p}{P_{opt}} \right) \quad (1)$$

$$G = R \left( \frac{hc}{q\lambda} \right) \frac{1}{\eta} \quad (2)$$

$$D^* = \frac{1}{NEP} \approx \sqrt{\frac{A}{2qI_d}} R. \quad (3)$$

Where  $I_p$  is the photo-excited current,  $P_{opt}$  the incident light power illuminated on the device,  $h$  the Planck's constant,  $c$  the velocity of the light,  $q$  the elemental charge,  $\lambda$  the incident light wavelength,  $\eta$  the quantum efficiency,  $A$  the effective area of the PD and  $I_d$  the dark current, respectively. Based on the above equations,  $R$  is calculated to be  $1.4 \times 10^2$  A/W (pure Bi<sub>2</sub>S<sub>3</sub> NR) and  $1.09 \times 10^3$  A/W (HGNs@Bi<sub>2</sub>S<sub>3</sub> NR) under the illumination of 560 nm light at a bias of 2 V.  $G$  is estimated to be  $2.68 \times 10^2$  (pure Bi<sub>2</sub>S<sub>3</sub> NR) and  $2.31 \times 10^3$  (HGNs@Bi<sub>2</sub>S<sub>3</sub> NR). What is more, by assuming  $\eta = 1$ ,  $D$  is determined to be  $2.45 \times 10^{12}$  and  $2.78 \times 10^{13}$  Jones or Bi<sub>2</sub>S<sub>3</sub> NRs and HGN@Bi<sub>2</sub>S<sub>3</sub> NRs, respectively. Table 1 summarizes the several important metrics of our devices and other Bi<sub>2</sub>S<sub>3</sub> nanostructures based photodetectors. Apparently, the key values of HGNs@Bi<sub>2</sub>S<sub>3</sub> NRs were not only higher than photodetectors based on pure Bi<sub>2</sub>S<sub>3</sub> NRs, but also than other Bi<sub>2</sub>S<sub>3</sub> nanostructures based devices [21]. This result suggests that the utilization of the plasmonic HGNs is a promising method for improving the optoelectronic devices.

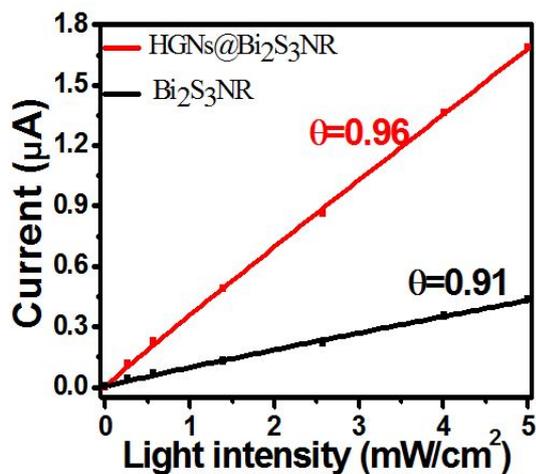
In addition to the high responsivity, conductive gain, and detectivity, the photocurrent of the plasmonic photodetector is highly dependent on the intensity of the excitation light. Figure 6 plots the photocurrents as a function of light intensity, from which one can find that the photocurrent increases almost linearly with increasing intensity. Careful fitting of both values finds that their relationship could be described by a power law [39]:  $I = AP^\theta$ , where  $I$ ,  $A$ ,  $P$  and  $\theta$  represent the photocurrent, the constant depending on the wavelength, the light power intensity and the exponent determining the response to light intensity, respectively. By fitting the curves, the  $\theta$  for HGN@Bi<sub>2</sub>S<sub>3</sub> NR is estimated to 0.96, slightly larger than that of pure Bi<sub>2</sub>S<sub>3</sub> NR (0.91). Such a relatively large exponent is probably due to the passivation effect of the ethylene glycol bis(3-mercaptopropionate) during surface modification process. These nearly integer exponents for both devices are much larger than that of device composed of Bi<sub>2</sub>S<sub>3</sub> nanocrystalline film (0.77),<sup>17</sup> which means that the density of trapping center in both devices is very low according to Rose's theory [40].



**Fig. 5.** (a) Schematic illustration of the HGNs@Bi<sub>2</sub>S<sub>3</sub> NR device. (b)  $I - V$  curves and responsivity of HGNs@Bi<sub>2</sub>S<sub>3</sub> NR in dark and under 560 nm light illumination, the inset is the TEM image of a single NR modified with HGNs. (c) Photoresponsive characteristics of the HGNs@Bi<sub>2</sub>S<sub>3</sub> NR device. (d) Summary of both dark current and photocurrent of 5 representative individual HGNs@Bi<sub>2</sub>S<sub>3</sub> NR based device.

**Table 1.** Comparison of the device performance of the present and other Bi<sub>2</sub>S<sub>3</sub> nanostructures based photodetectors.

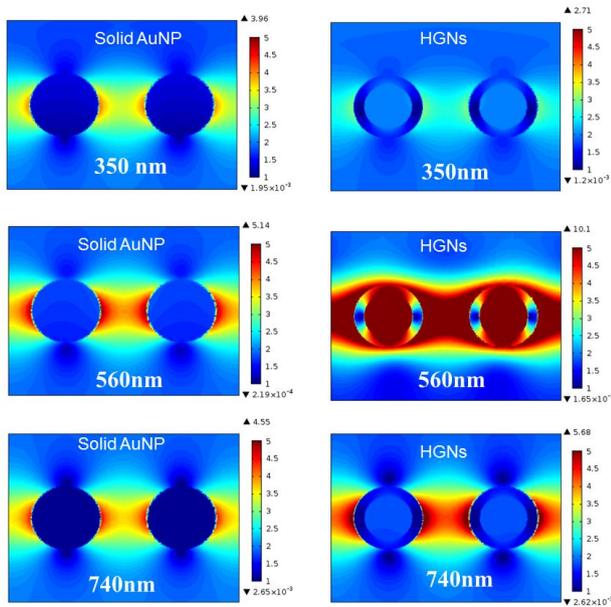
Materials	On/off ratio	$R/AW^{-1}$	$G$	$D^*/cmHz^{1/2}W^{-1}$	Ref.
HGNs@Bi <sub>2</sub> S <sub>3</sub> NR	1066	$1.09 \times 10^3$	$2.31 \times 10^3$	$2.78 \times 10^{13}$	our work
Bi <sub>2</sub> S <sub>3</sub> NR	147	$1.4 \times 10^2$	$2.68 \times 10^2$	$2.45 \times 10^{12}$	Our work
Bi <sub>2</sub> S <sub>3</sub> crystalline	—	1.7	~40	~ $10^{11}$	[21]
Bi <sub>2</sub> S <sub>3</sub> NW	3.4	—	—	—	[38]



**Fig. 6.** The fitting of the relationship between the photocurrent and light intensity for devices based on both pure Bi<sub>2</sub>S<sub>3</sub> NR and HGNs@Bi<sub>2</sub>S<sub>3</sub> NR.

The large enhancement of HGN@Bi<sub>2</sub>S<sub>3</sub> NRs may be related to the strong LSPR, which is verified by FEM simulation. Figure 7 illustrates the electric field intensity distribution of the HGNs under light illumination with wavelengths of 350, 560 and 740 nm. Obviously, when shined by 560 nm light illumination, the HGNs can induce large-area hot spots with high field intensity. The peak dipolar electric energy according to the color bar, is more than 10-fold, which is much stronger than the ~2.7-fold (350 nm) and 5.6-fold (740 nm), respectively. In addition, the field enhancement of HGNs under 560 nm light is much stronger than HGNs with the same diameter (5.1-fold). Such superb optical properties can be ascribed to their unique geometry: unlike the conventional solid gold nanoparticles with only one exterior surface, the HGNs have one more surface within the cavity. The two faces are able to facilitate coupling at light illumination, thus leading to relatively large field enhancement [41], which is highly beneficial for light detection. Without question, compared with traditional plasmonic photodetectors using solid no-

ble metal nanoparticles, the utilization of HGNs for enhancing nano-photodetector is advantageous in the following two aspects. (1) Large enhancement at low cost. Although the HGNs consume relatively few gold in comparison with SGNs, their field enhancement in contrast is much stronger. (2) Tunable LSPR band. As discussed above, the LSPR of HGNs can be easily adjusted in a range from 500–610 nm by tuning the aspect ratio. This convenience will make it possible to develop ultrasensitive nano-photodetectors with tunable spectral selectivity.



**Fig. 7.** The electrical field distribution of both solid AuNP and HGNs under light illumination of 350, 560, and 740 nm. The solid AuNP has a diameter of 30 nm, and the HGN has a diameter and aspect ratio of 30 nm and 70%, respectively.

When the HGNs with strong LSPR effect were modified on the surface of  $\text{Bi}_2\text{S}_3$  nanoribbons, the photocurrent, responsivity, and detectivity can be enhanced by the following three mechanisms: (1) Increased capability of light absorption by field enhancement. As a result, more photo-generated electron-hole pairs were generated, leading to enhanced photocurrent. (2) Accelerated separation of electron-hole pairs owing to the enhanced electromagnetic field surrounding the plasmonic HGNs [42], [43], [44]. (3) Direct electron transfer (DET). When irradiated by light illumination, the energetic hot electrons with energy as high as 1–4 eV will directly inject into the conduction band of the nearby semiconductor nanostructures [45], [46]. As a result, the photocurrent and responsivity were considerably increased.

## 4 Conclusion

In summary, we have demonstrated a simple strategy to improve the performance of nano-photodetector. The assembled HGN@ $\text{Bi}_2\text{S}_3$  NR photodetector is highly sensitive to 560 nm illumination light with good reproducibility. Moreover, the HGNs decorated photodetector exhibits enhanced responsivity, conductivity gain and detectivity in comparison with pure  $\text{Bi}_2\text{S}_3$  NR and other  $\text{Bi}_2\text{S}_3$  monocrystalline based devices. Theoretical simulation based FEM reveals that the enhancement in device performance is related to the LSPR effect of the HGNs, which may increase the light absorption of the HGNs@ $\text{Bi}_2\text{S}_3$  NR heterojunction. This study suggests that plasmonic HGNs are highly reliable and efficient platforms to boost the device performance of 1-D semiconductor nanostructures based optoelectronic devices.

## 5 Methods

Synthesis and structural analysis of  $\text{Bi}_2\text{S}_3$  NRs and hollow gold nanoparticles: The  $\text{Bi}_2\text{S}_3$  NRs were synthesized *via* a simple solvothermal approach. Briefly, NaOH,  $\text{Bi}(\text{NO}_3)_3$ , and  $\text{Na}_2\text{S}_2\text{O}_3$  were added sequentially into a mixed solution of glycerol and water (2:1 by volume), leading to a solution with concentration of 50 mM of  $\text{Bi}^{3+}$ , 2 M of NaOH, and 100 mM of  $\text{Na}_2\text{S}_2\text{O}_3$ . The mixed solution was then transferred into a Teflon-lined autoclave (50 mL capacity) which was sealed and maintained at 160°C. After keeping for 20 h, dark-gray precipitate floating on the top of the solution was observed. The product was centrifuged and washed with distilled water for several times to remove possible contaminant and ions. The HGNs were prepared by a sacrificial galvanic replacement reaction between Co nanoparticles and aqueous  $\text{HAuCl}_4$ . The Co nanoparticles were firstly synthesized by the following process: 1 M sodium borohydride solution (100  $\mu\text{L}$ ), 0.1 M trisodium citrate dehydrate (400  $\mu\text{L}$ ), 0.4 M cobalt hexahydrate solution (100  $\mu\text{L}$ ) and ultra-pure water (100 mL) were put into a three-neck flask and stirred to form a homogenous solution. To avoid the oxidation of the Co, the reaction solution was purged with  $\text{N}_2$  gas flow during the whole reaction process. After the formation of Co nanoparticles, 50  $\mu\text{L}$  0.1 M of chloroauric acid trihydrate solution was dropped into the above solution. Finally, the HGN was collected by centrifuging the HGN solution at 8000 rpm for 20 min. The phase of the  $\text{Bi}_2\text{S}_3$  NR was examined using an X-ray diffraction (XRD, Rigaku D/Max-yB with Cu  $K\alpha$  radiation). The morphology and crystal

structure were carried out on a SIRION 200 FEG scanning electron microscopy (FESEM), transmission electron microscopy (TEM), and high-resolution transmission electron microscopy (HRTEM, JEOL JEM-2010, at 200 kV). The optical absorption properties were performed on a UV-Vis absorption spectrophotometry (CARY 5000). The chemical composition of the product was analyzed using energy dispersive X-ray spectrometer (EDS), and X-ray photoelectron spectrometer (XPS).

**Device construction and characterization:** The Bi<sub>2</sub>S<sub>3</sub> NRs based device was fabricated *via* the following procedures. Firstly, the as-synthesized Bi<sub>2</sub>S<sub>3</sub> nanoribbons were dispersed on the SiO<sub>2</sub> (300 nm)/Si substrate with a desired density. Then two parallel In electrodes were defined on the NR through a photolithography and e-beam evaporation method. Finally, to further decorate with hollow gold nanoparticles, the as-assembled device was immersed into an ethylene glycol bis(3-mercaptopropionate) solution for 1 h and HGN aqueous solution for 12 h sequentially. The optoelectronic properties of the fabricated photodetectors were measured on a Keithley 4200-SCS I-V system equipped with a monochromatic light source system (Princeton, ARC-SP-2155).

**Theoretical simulation:** The theoretical modeling was carried out by using a finite element method (FEM) to study the unit cell. A periodic boundary condition is used along x direction and transverse magnetic (TM) polarized light was assumed to be normally incident from the top of the structure. The calculated region is bordered by perfectly matched layers (PML) to reduce the influence of light reflection. The optical data of Au were from Sopra S. A. company database.

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